



Statistics for Astronomers: Lecture 15, 2019.04.22

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Recall: Priors

(from Ivezić et al.)

In terms of information, priors can be **informative** or **“non-informative”**.

Informative prior

Specific information about parameter(s). Progressively increasing amounts of data \implies posterior is evidence-dominated.

Example: “Data from the past ten years suggests that there is a 2% change of rain in Morelia today between 2 and 3 PM.”

Non-informative prior

Vague information about parameters, typically based on general principles/objective information (also called objective prior). “Light” modification to observations \implies posterior is likelihood-dominated.

Example: “The flux from this star is non-negative” ($0 \leq F < \infty$).

This is also an example of an **improper prior**, as it does not integrate to unity.

However, we are still OK if the resulting posterior is well-defined

(bus example from last week – $p(\tau|I) \propto 1/\tau, t \leq \tau < \infty$).

The Principle of Indifference is a classic example of an uninformative prior.



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Recall: Bayesian point/location and interval estimates

Once $p(\theta|\text{data})$ is computed, we can compute the location estimates (mean, median, mode).

For example, the **Bayesian estimator** of the parameter mean is $\bar{\theta} = \int d\theta \theta p(\theta|\text{data})$.

We can also compute Bayesian interval estimates, also called **posterior intervals** or **credible intervals** (abbreviated in these lectures as CrI).

One example of a $100(1 - \alpha)\%$ CrI is $[a, b]$ such that

$$\int_{-\infty}^a d\theta p(\theta|\text{data}) = \int_b^{\infty} d\theta p(\theta|\text{data}) = \alpha/2.$$

Another type of CrI is the **highest posterior density** (HPD) interval, defined as the **narrowest interval** that contains $100(1 - \alpha)\%$ of the posterior probability.



Recall: Numerical computation of HPD interval

- 1 Obtain N random deviates $x[i]$ drawn from the posterior density distribution.
- 2 Sort them in ascending order.
- 3 For each $x[i]$, find the point that is $w = (1 - \alpha)N$ points away.
- 4 Compute the widths $w[i] = x[w + i] - x[i]$.
- 5 Find the location $i = i_0$ corresponding to the smallest width. The HPD interval is then $(x[i_0], x[i_0 + w])$.

Write your own script! You'll need it for your research if you're using Bayesian inference.



Recall: Prior-dominated posterior

Prior: $p(S) \propto S^{-5/2}$.

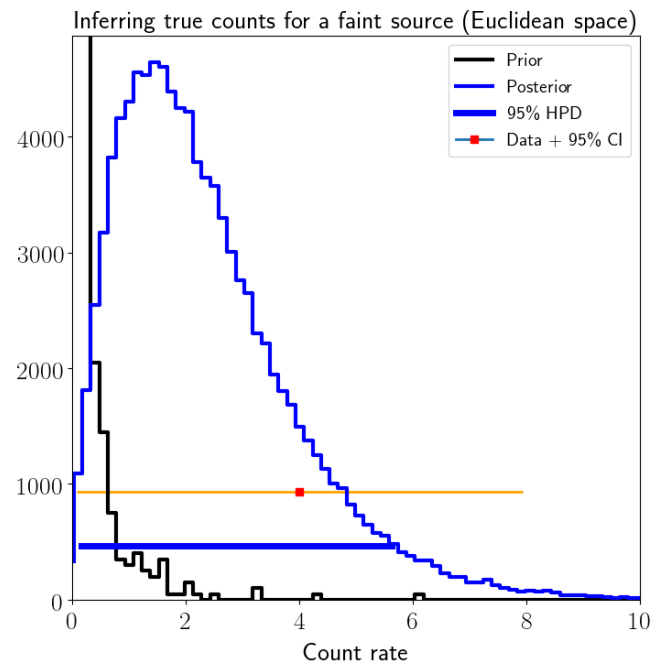
Likelihood of obtaining data $S_{\text{obs}} = 4$ from Poissonian uncertainties acting on S :

$$\mathcal{L}(S) \propto S^{S_{\text{obs}}} \exp[-S] = S^4 \exp[-S].$$

$$\text{Posterior } p(S|S_{\text{obs}}) \propto S^{3/2} \exp[-S]$$

$$= \text{Gamma}\left(\frac{5}{2}, 1\right).$$

⇒ Mean: $5/2$; Mode: $3/2$. Can also compute HPD (homework).



Example of prior-dominated posterior

from Andreon, "Bayesian Methods for the Physical Sciences".

Andreon et al. 2009 mass measurement for most distant ($z \geq 2$) galaxy cluster, JKCS041.

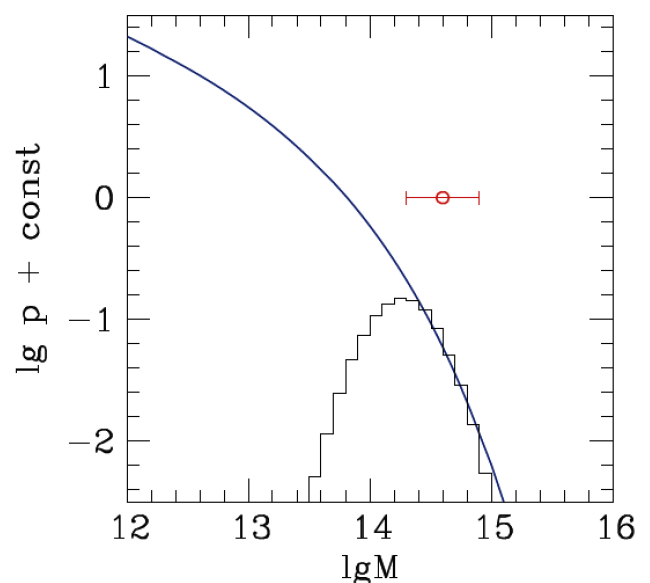
Mass estimate important to constrain parameters of Λ CDM model.

Observation: $\log M/M_{\odot} = 14.6 \pm 0.3$.

Prior: Schechter mass function.

Prior changes drastically near observed value, similar to previous example.

Posterior mean is therefore lower than observed value: $\log M/M_{\odot} = 14.3 \pm 0.3$.
(lower by 2x!)



Recall: Maximum likelihood and Fisher information

The variance associated with the MLE estimate is bounded below by the reciprocal of the Fisher information (Cramér-Rao Bound).

The square-root of the reciprocal of the Fisher information is therefore a lower bound to the standard deviation of the MLE estimate.

Given the likelihood $\mathcal{L}(\theta)$, the Fisher information is given by

$$\mathcal{I}(\theta) = \mathbb{E} \left[\left(\frac{\partial \ln \mathcal{L}}{\partial \theta} \right)^2 \right] \quad (\text{under some regularity conditions}) = -\mathbb{E} \left[\frac{\partial^2 \ln \mathcal{L}}{\partial \theta^2} \right]$$

The Fisher information is related to the [curvature](#) of the log-likelihood near the MLE value.



The Jeffreys prior

One example of a non-informative prior that is also [invariant](#) over transformation of the random variable (the form of the dependence on the variable doesn't change).

$\pi_J(\theta) \propto \sqrt{\mathcal{I}(\theta)}$, where $\mathcal{I}(\theta)$ is the Fisher information.

In the multidimensional case, replace \mathcal{I} above with the determinant of the Fisher information matrix.

As previously noted, the Fisher information is related to the variance. [This form of prior is ideal for scale parameters.](#)

Invariance: Let ψ be some function of θ (e.g., if θ is the probability of a coin flip resulting in a head, then $\psi = \frac{\theta}{1-\theta}$, the [odds ratio](#), is a function of θ). We then have

$$\begin{aligned} \pi_J(\psi) &= \pi_J(\theta) \left| \frac{d\theta}{d\psi} \right| \propto \sqrt{\mathcal{I}(\theta) \left(\frac{d\theta}{d\psi} \right)^2} = \sqrt{\mathbb{E} \left[\left(\frac{\partial \ln \mathcal{L}}{\partial \theta} \right)^2 \right] \left(\frac{d\theta}{d\psi} \right)^2} = \sqrt{\mathbb{E} \left[\left(\frac{d\theta}{d\psi} \frac{\partial \ln \mathcal{L}}{\partial \theta} \right)^2 \right]} \\ &= \sqrt{\mathbb{E} \left[\left(\frac{\partial \ln \mathcal{L}}{\partial \psi} \right)^2 \right]} = \sqrt{\mathcal{I}(\psi)}. \end{aligned}$$

The form of the dependence on the parameter is the same, regardless of whether it is θ or ψ .



Jeffreys prior example: a Bernoulli trial

Let θ be the probability of “success” in a Bernoulli trial. We perform one trial and obtain a value $X = x$.

The likelihood associated with this observation is $\mathcal{L}(\theta) \propto \theta^x(1-\theta)^{1-x} = \text{Beta}(x+1, 2-x)$
 $\implies \ln \mathcal{L} = x \ln \theta + (1-x) \ln(1-\theta) \implies \frac{\partial \ln \mathcal{L}}{\partial \theta} = \frac{x}{\theta} - \frac{1-x}{1-\theta} = \frac{x-\theta}{\theta(1-\theta)}$.

$$\mathcal{I}(\theta) = \mathbb{E} \left[\left(\frac{\partial \ln \mathcal{L}}{\partial \theta} \right)^2 \right] = \frac{1}{\theta(1-\theta)}$$

$\implies \pi_J(\theta) \propto \frac{1}{\sqrt{\theta(1-\theta)}} = \text{Beta}\left(\frac{1}{2}, \frac{1}{2}\right)$; prior mean: $\frac{1/2}{1/2+1/2} = 0.5$ as expected.

Posterior: $p(\theta|\text{data}) \propto \mathcal{L}(\theta)\pi_J(\theta) = \text{Beta}(x+1, 2-x) \times \text{Beta}\left(\frac{1}{2}, \frac{1}{2}\right) = \text{Beta}\left(x + \frac{1}{2}, \frac{3}{2} - x\right)$.
 Posterior mean: $0.5(x + 0.5) = 0.5$ (sample mean) + 0.5 (prior mean). Effective sample size: 2.

Note that the posterior and prior are both Beta distributions. In such a case, we say that the Beta distribution is the **conjugate prior** to a Bernoulli likelihood. The Beta distribution is also conjugate to binomial likelihoods.



Jeffreys priors for a univariate normal distribution

$$\mathcal{L}(x|\mu, \sigma) = \frac{1}{\sqrt{2\pi\sigma^2}} \exp \left[-\frac{1}{2} \left(\frac{x-\mu}{\sigma} \right)^2 \right] \implies \ln \mathcal{L} = -\ln \sigma - \frac{1}{2} \left(\frac{x-\mu}{\sigma} \right)^2.$$

$$\frac{\partial \ln \mathcal{L}}{\partial \mu} = \left(\frac{x-\mu}{\sigma^2} \right).$$

$$\mathbb{E} \left[\left(\frac{\partial \ln \mathcal{L}}{\partial \mu} \right)^2 \right] = \mathbb{E} \left[\left(\frac{x-\mu}{\sigma^2} \right)^2 \right] = \frac{1}{\sigma^2} \mathbb{E} \left[\left(\frac{x-\mu}{\sigma} \right)^2 \right] = \frac{1}{\sigma^2} \mathbb{E}[z] = \frac{1}{\sigma^2} \propto \text{constant}.$$

\implies uniform prior for μ .

$$\frac{\partial \ln \mathcal{L}}{\partial \sigma} = -\frac{1}{\sigma} + \frac{(x-\mu)^2}{\sigma^3}; \quad \frac{\partial^2 \ln \mathcal{L}}{\partial \sigma^2} = \frac{1}{\sigma^2} - 3 \frac{(x-\mu)^2}{\sigma^4}.$$

$$\mathbb{E} \left[-\frac{\partial^2 \ln \mathcal{L}}{\partial \sigma^2} \right] = \frac{1}{\sigma^4} \mathbb{E} \left[3 \left(\frac{x-\mu}{\sigma} \right)^2 - 1 \right] \propto \frac{1}{\sigma^2}.$$

Therefore, the prior for σ is $\pi_J(\sigma) \propto \frac{1}{\sigma}$

\implies logarithmic prior for σ .



More on priors

See <https://bit.ly/2KW92Pt> for a good discussion of the applicability of this procedure to problems in fundamental physics.

